

Strong coupling between a single artificial atom and an integrated silicon carbide microresonator

Daniil M. Lukin^{1†}, Dominic Catanzaro¹, Melissa A. Guidry¹, Eran Lustig¹, Misagh Ghezellou², Joshua Yang¹, Hiroshi Abe³, Takeshi Ohshima^{3,4}, Jawad Ul-Hassan², Jelena Vučković¹

¹*E. L. Ginzton Laboratory, Stanford University, Stanford, California 94305, USA*

²*Department of Physics, Chemistry and Biology, Linköping University, SE-58183, Linköping, Sweden*

³*National Institutes for Quantum Science and Technology, Takasaki, Gunma 370-1292, Japan*

⁴*Department of Materials Science, Tohoku University, Sendai 980-8579, Japan*

[†]*dlukin@stanford.edu*

Abstract: The strong coupling regime between a photonic cavity and an artificial atom in 4H-Silicon Carbide-on-insulator photonics is demonstrated, using a high-finesse whispering gallery mode resonator and a single silicon vacancy center.

© 2024 The Author(s)

An optically-addressable solid-state spin strongly coupled to a photonic resonator enables deterministic gates between a flying qubit (a photon) and a quantum memory qubit (electron spin), a valuable resource for building scalable photon-mediated entanglement [1, 2]. Whereas emitter-cavity systems with large cooperativity but in the weak coupling regime enable the modification of cavity transmission spectrum [3] for efficient read-out of the spin state, in a strongly-coupled system, an entangled light-matter state can be directly produced. A large cooperativity requires that the atom-cavity coupling rate g^2 far exceeds the product of the cavity decay rate κ and the atom dephasing/decay rates (γ). Therefore, a system can still have a large cooperativity even when $g < \kappa$, *i.e.* when cavity loss rate dominates. In contrast, for the strong coupling metric, the rates κ and γ are treated on an equal footing. A system is considered strongly coupled if the condition $4g > \kappa, \gamma$ is met. The condition $4g > \kappa$ is challenging to satisfy in integrated photonics and is typically the limiting factor for achieving strong coupling. This is especially true for color centers, which have a relatively small dipole moment compared to, for instance, semiconductor quantum dots.

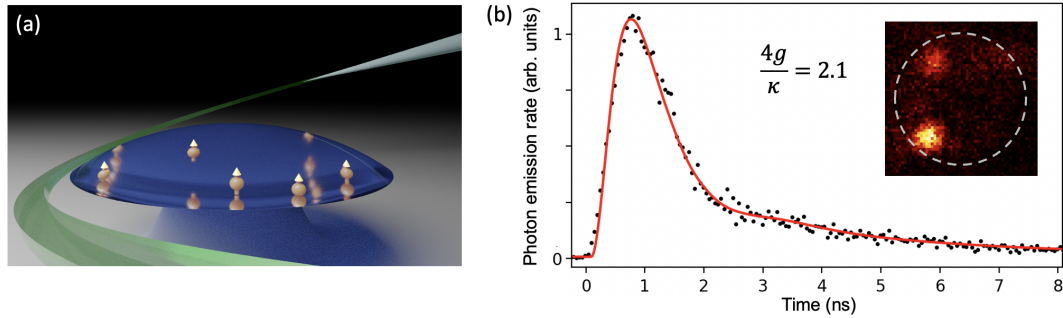


Fig. 1. (a) Illustration of photonic device. A low-roughness resonator is fabricated via photoresist reflow pattern transfer [4], and addressed via waveguide-fiber interface [5]. The resonator contains multiple randomly-placed emitters. (b) By spectrally tuning the resonator to the edge of the emitter inhomogeneous distribution, we isolate a detuning where only a single emitter is well-coupled to the cavity (spatial distribution scan shown in inset). By fitting the measured temporal dynamics of the emitter cavity system after excitation with a pulsed laser, we obtain $4g/\kappa = 2.1$.

Here, we integrate silicon vacancy color centers [6, 7] into high-finesse whispering gallery mode resonators in 4H-SiC on Insulator [8], and reach the strong coupling regime through a dramatic reduction of cavity loss rate κ . Following the technique described in Refs. [4, 5, 9], we fabricate whispering gallery mode resonators with finesse exceeding 10^3 . An illustration of the device is shown in Fig. 1(a). To observe strong coupling between a single emitter and the resonator, we tune the optical mode via cryogenic gas condensation to the edge of the inhomogeneous distribution of the emitters, such that color center emission into the cavity mode is dominated by a single emitter, as confirmed via second-order autocorrelation measurements as well as a photoluminescence map of the

device. The optical mode loaded quality factor of $8.55 \cdot 10^5$ is measured via laser transmission, corresponding to $\kappa = 383$ MHz. By exciting the emitter via a pulsed above-resonant mode-locked laser, we observe fast temporal dynamics of photon emission from the cavity mode (Fig. 1(b)). Under above-resonant excitation, the emitter undergoes spectral diffusion, which is modeled by sampling dynamics for a distribution of emitter detunings. The result of the model captures the experimentally observed dynamics with the emitter-cavity coupling strength as a free parameter. The obtained coupling rate $g = 198$ MHz, higher compared to the previous demonstration [9], is consistent with the reduced mode volume of the device. Hence, $4g/\kappa = 2.1$, which exceeds the threshold for strong coupling by two times. We then characterize the single-photon purity and indistinguishability of the light emitted by the system, observing $g^{(2)}(0) = 0.02$ and interference visibility of 0.75, which confirms that the realized strongly coupled system is a source of pure and indistinguishable single photons.

This demonstration is a step toward deterministic gates between stationary and flying qubits mediated by the solid-state artificial atom spin-optical interface.

Acknowledgements

This work was supported by the Vannevar Bush Faculty Fellowship. J.U.H. acknowledges support from Swedish Research Council (grant No. 2020-05444), Knut and Alice Wallenberg Foundation (grant No. KAW 2018-0071), and the EU H2020 project QuanTELCO (grant No. 862721). T.O. acknowledges grants JSPS KAKENHI 20H00355 and 21H04553. Part of this work was performed at the Stanford Nanofabrication Facility (SNF) and the Stanford Nano Shared Facilities (SNSF).

References

1. H. J. Kimble, "The quantum internet," *Nature* **453**, 1023–1030 (2008).
2. I. Fushman, D. Englund, A. Faraon, N. Stoltz, P. Petroff, and J. Vuckovic, "Controlled phase shifts with a single quantum dot," *Science* **320**, 769–772 (2008).
3. E. Waks and J. Vuckovic, "Dipole induced transparency in drop-filter cavity-waveguide systems," *Phys. review letters* **96**, 153601 (2006).
4. D. M. Lukin, M. A. Guidry, M. Ghezellou, D. Catanzaro, J. Yang, H. Abe, T. Ohshima, J. Ul-Hassan, and J. Vučković, "Multimode cavity quantum electrodynamics in 4H-silicon carbide-on-insulator photonics," in *CLEO: Fundamental Science*, (Optica Publishing Group, 2023), pp. FTu3C–4.
5. D. Catanzaro, D. M. Lukin, E. Lustig, M. A. Guidry, and J. Vučković, "Cryogenic fiber-coupled waveguide probe co-integrated with electrical control lines," in *CLEO: Science and Innovations*, (Optica Publishing Group, 2023), pp. JTu2A–47.
6. R. Nagy, M. Niethammer, M. Widmann, Y.-C. Chen, P. Udvarhelyi, C. Bonato, J. U. Hassan, R. Karhu, I. G. Ivanov, N. T. Son *et al.*, "High-fidelity spin and optical control of single silicon-vacancy centres in silicon carbide," *Nat. communications* **10**, 1–8 (2019).
7. D. Liu, F. Kaiser, V. Bushmakina, E. Hesselmeier, T. Steidl, T. Ohshima, N. T. Son, J. Ul-Hassan, Ö. O. Soykal, and J. Wrachtrup, "The silicon vacancy centers in sic: determination of intrinsic spin dynamics for integrated quantum photonics," *arXiv preprint arXiv:2307.13648* (2023).
8. D. M. Lukin, C. Dory, M. A. Guidry, K. Y. Yang, S. D. Mishra, R. Trivedi, M. Radulaski, S. Sun, D. Vercruyse, G. H. Ahn *et al.*, "4H-silicon-carbide-on-insulator for integrated quantum and nonlinear photonics," *Nat. Photonics* **14**, 330–334 (2020).
9. D. M. Lukin, M. A. Guidry, J. Yang, M. Ghezellou, S. D. Mishra, H. Abe, T. Ohshima, J. Ul-Hassan, and J. Vučković, "Two-emitter multimode cavity quantum electrodynamics in thin-film silicon carbide photonics," *Phys. Rev. X* **13**, 011005 (2023).